



# The Effects of Heat Generation/Absorption and Magnetic Field on Marangoni Flow of a Nanofluid Containing Gyrotactic Microorganisms Over a Stretching Surface

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## Abstract

The current study examines the numerical approximation of Marangoni flow of a Nanofluid containing gyrotactic microorganisms over a stretching surface in the presence heat generation and magnetic field. The governing equations are written down and converted into a system of nonlinear ordinary differential equations. The resultant system is solved numerically. Discussion and analysis are conducted on the impacts of the factors on the velocity, temperature, concentration, and density of motile microbe. Additionally, computed and explored for various embedded factors in the issue statements are the density of the motile microbe, Nusselt number, Sherwood number, and skin friction coefficient.

**Keywords:** Marangoni Flow; Stretching Surface; Heat, And Mass Transfer; Gyrotactic Micro-Organism; Magnetic Field; Heat Generation

## Introduction

The mass transfer across an interface between two fluids caused by a differential in surface tension is known as the Gibbs-Marangoni effect. This aspect is referred to as thermo-capillary convection or Bénard-Marangoni convection in the event of temperature dependency. In fact, Marangoni convection, which is brought on by changes in surface tension gradients (i.e., temperature and concentration gradients), finds significant use in the processing of silicon wafers, soap films, nucleation of vapor bubbles, spreading of thin films, and other semi-conductor-related fields. In their study of Marangoni-driven boundary layer flow in nanofluids, Nazar et al. [1] found that heat transmission is enhanced more effectively by nanoparticles with poor thermal conductivity. Marangoni convection has been studied by several scholars in a variety of geometries [2-5]. Arifin et al. [6] inspected the influence of radiation on Marangoni convection boundary layer flow of nanofluid and found that a raise in thermal radiation parameter reduces the quota of heat transmitted at the surface. Copper-water nanofluid flow and heat transport through a porous medium disc were studied by Yanhai and Liancun [7]. Ambrish et al. [8] study the mathematical Marangoni Convection MHD Flow of Carbon Nanotubes across a Porous Medium. The impact of MHD Marangoni Convective Flow of Nanofluid through a porous medium with heat and mass transfer was investigated by Venkateswara

Raju et al. [9]. The habitation of heavy microbe collecting on top of lighter water is what causes the phenomena of bio-convection in nanofluid, which is the subject of current investigation. The establishment of the bio-convection process within the system happens when the heavier bacteria fall into the water and are replaced by up floating microbe. A macroscopic motion is caused by the mesoscale phenomena of moving motile Microbe (convection). The literature: [10-21] covers the evolution of theoretical work in bio-convection. The properties of heat transport in the laminar boundary layer of a viscous fluid over a stretched sheet with viscous dissipation or frictional heating and internal heat generation were examined by Vajravelu and Haidilao [22]. Vajravelu and Hadjinolaou [22] estimated the volumetric rate of heat production in their investigation to be

$$q^* = \begin{cases} Q_0(T - T_\infty) & \text{for } T \geq T_\infty \\ & \text{for } T \leq T_\infty \end{cases}$$

Where  $Q_0$  is the heat generation/absorption constant. The relationship mentioned above, as Vajravelu and Hadjinolaou [22] stated, is accurate as a state approximation for various exothermic processes using  $T_\infty$  as the onset temperature. They employed

$q^{*} = Q_0(T - T_{\infty})$  when the input temperature was not less than  $T_{\infty}$ . It is clear from the research reviewed above that the issue of Marangoni flow, heat, mass, and motile microbe transmission over a stretching surface has not yet been taken into consideration.

**Mathematical Model**

As exhibited in Figure 1, the current study takes the effects of heat production into account as a constant two-dimensional incompressible Marangoni convective boundary layer flow of a nanofluid across a stretched surface. The coordinates are  $x$

and  $y$  along the stretching surface and normal to it, respectively. Suppose that flow occurs across  $y \geq 0$ . A uniform transverse magnetic field of strength  $B_0$  is implemented in  $y$ -direction, The magnetic Reynolds number is taken to be small enough so that the induced magnetic field can be neglected. Also assume that  $T(x)$ ,  $C(x)$  and  $N(x)$  represent the fluid's temperature, concentration, and motile microbe density, respectively. Let's assume that  $T_{\infty}$ ,  $C_{\infty}$ , and  $N_{\infty}$  are the corresponding amounts in the free stream. The surface tension depends on temperature, concentration, and density of motile Microbe such that

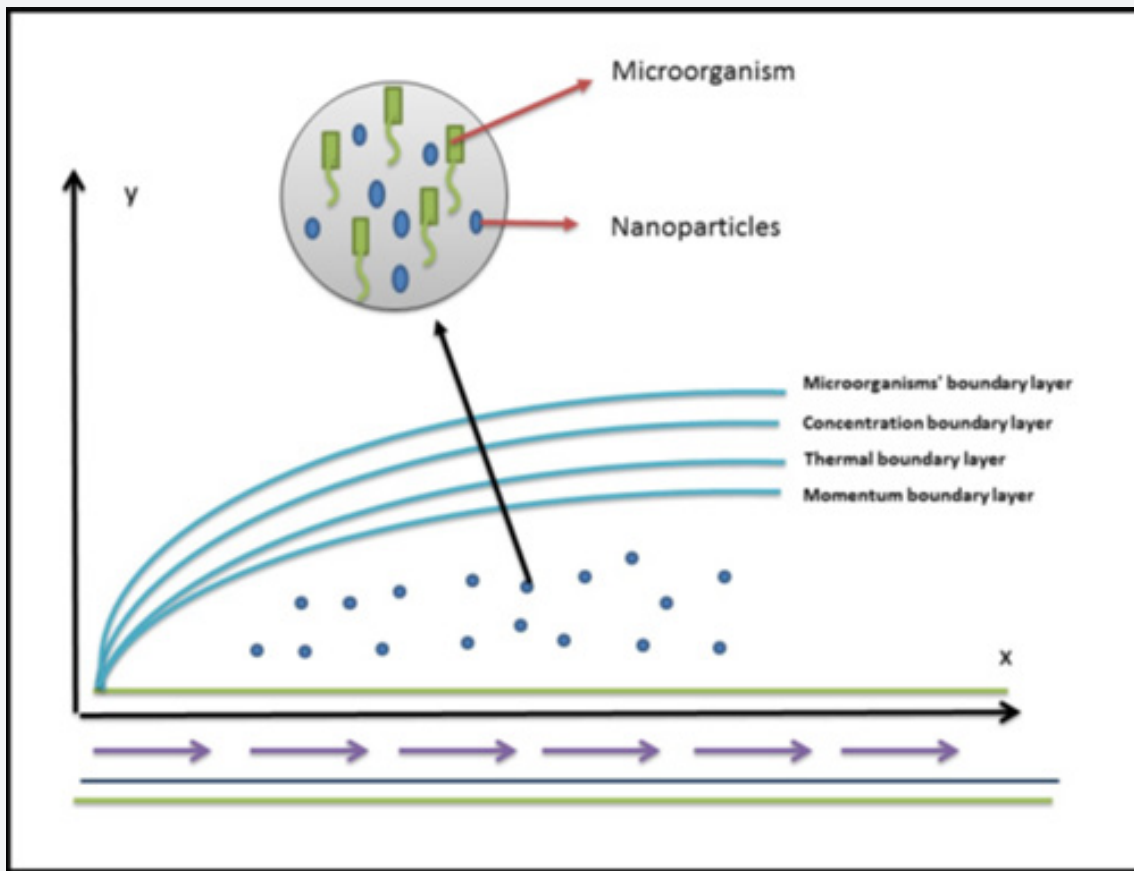


Figure 1: Schematic of flow above stretching surface.

$$\sigma = \sigma_0[1 - \gamma(T - T_{\infty}) - \gamma^*(C - C_{\infty}) - \gamma^{**}(N - N_{\infty})]$$

Where  $\sigma_0$  is interface surface tension ratio and

$$\gamma = -\frac{1}{\sigma_0} \left( \frac{\partial \sigma}{\partial T} \right)_{C,N}, \quad \gamma^* = -\frac{1}{\sigma_0} \left( \frac{\partial \sigma}{\partial C} \right)_{T,N}, \quad \gamma^{**} = -\frac{1}{\sigma_0} \left( \frac{\partial \sigma}{\partial N} \right)_{T,C}$$

The following are the governing equations for this investigation: (Figure 1).

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \tag{1}$$

$$\left( u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right) = \nu \left( \frac{\partial^2 u}{\partial y^2} \right) - \frac{\sigma^* B_0^2}{\rho_{\infty}} u, \tag{2}$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha \left( \frac{\partial^2 T}{\partial y^2} \right) + \tau \left[ D_B \left( \frac{\partial T}{\partial y} \frac{\partial C}{\partial y} \right) + \frac{D_T}{T_{\infty}} \left( \frac{\partial T}{\partial y} \right)^2 \right] + Q_0(T - T_{\infty}) \tag{3}$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_B \left( \frac{\partial^2 C}{\partial y^2} \right) + \frac{D_T}{T_{\infty}} \left( \frac{\partial^2 T}{\partial y^2} \right), \tag{4}$$

$$\frac{\partial N}{\partial x} + v \frac{\partial N}{\partial y} = D_n \left( \frac{\partial^2 N}{\partial y^2} \right) - \frac{bW_c}{C_w - C_\infty} \left[ \frac{\partial}{\partial y} N \left( \frac{\partial C}{\partial y} \right) \right], \quad (5)$$

with boundary conditions

$$\text{at } y = 0: \mu \frac{\partial u}{\partial y} = \frac{\partial \sigma}{\partial x}, \quad v = 0, \quad T = T_w, \quad C = C_w, \quad N = N_w$$

$$\text{as } y \rightarrow \infty: u \rightarrow 0, \quad T \rightarrow T_\infty, \quad C \rightarrow C_\infty, \quad N \rightarrow N_\infty \quad (6)$$

where  $\nu$  is the kinematic viscosity, where  $\sigma^*$  is the electrical conductivity,  $\rho_\infty$  is the density of fluid,  $\tau$  the ratio of the effective heat capacitance of the nanoparticle to that of the base fluid,  $\alpha$  is the thermal diffusivity of the fluid,  $T_w$  is the surface temperature,  $C_w$  is the surface concentration,  $N_w$  is the surface density of microorganism,  $D_B$ ,  $D_T$  and  $D_n$  are the Brownian diffusion coefficient, thermophoresis diffusion coefficient and diffusivity of micro-organisms, respectively,  $b$  is the chemotaxis constant,  $W_c$  is maximum cell swimming speed,  $bW_c$  is assumed to be constant, and  $Q_0$  is the heat generation or absorption coefficient such that  $Q_0 > 0$  corresponds to heat generation while  $Q_0 < 0$  corresponds to heat absorption.

Select a stream function  $\psi(x, y)$  such that the equation of continuity is fulfilled. So

$$u = \frac{\partial \psi}{\partial y}, \quad v = -\frac{\partial \psi}{\partial x}.$$

Introducing the following similarity variables:

$$\theta(\gamma, z) = \nu X f(\xi), \quad \xi = \frac{y}{L}, \quad X = \frac{x}{L}, \quad \xi(\xi) = \frac{T - T_\infty}{T_w - T_\infty}, \quad T_w - T_\infty = T_0 X^2,$$

$$\phi = \frac{C - C_\infty}{C_w - C_\infty}, \quad C_w - C_\infty = C_0 X^2, \quad X = \frac{N - N_\infty}{N_w - N_\infty}, \quad N_w - N_\infty = N_0 X^2,$$

Where  $T_0$  is the reference temperature,  $C_0$  is the reference concentration and  $N_0$  is the reference density of micro-organisms. Equations (2), (3), (4) and (5) can be written as

$$f''' - f'^2 + ff'' - \xi f' = 0, \quad (7)$$

$$\theta'' + \text{Pr} f \theta' - 2 \text{Pr} f' \theta + \text{Pr} N_b \phi \theta' + \text{Pr} N_t \theta'^2 + \lambda \text{Pr} \theta = 0, \quad (8)$$

$$\phi'' = \text{Sc} f \phi' + \frac{N_t}{N_b} \theta'' = 0, \quad (9)$$

$$X'' + \text{Le} f X' - 2 \text{Le} f' X - \text{Pe} [(x + \delta) \phi'' + \phi' X'] = 0, \quad (10)$$

with boundary conditions

$$\eta = 0: f = 0, \quad f''(0) = -2(1 + \delta^*), \quad \theta = 1, \quad \phi = 1, \quad X = 1$$

$$\eta \rightarrow 0: f' = 0, \quad \theta = 0, \quad \phi = 0, \quad X = 0 \quad (11)$$

$$\xi = \frac{L^2 \sigma^* B_0^2}{\rho_\infty}$$

where  $\rho_\infty$  is the magnetic field parameter,  $\text{Pr} = \frac{\nu}{\alpha}$  is the Prandtl number,  $N_b = \tau D_B (C_w - C_\infty) / \nu$  is the Brownian motion parameter,  $N_t = \tau D_T (T_w - T_\infty) / \nu T_\infty$  is the thermophoresis parameter,  $\lambda = Q_0 L^2 / \alpha$  is the heat-source/sink parameter,  $\text{Sc} = \nu / D_B$  is the Schmidt number,  $\text{Le} = \nu / D_n$  is the Lewis number,  $\text{Pe} = bW_c / D_n$  is the Peclet number,  $\delta = N_\infty / (N_w - N_\infty)$  is motile Microbe parameter,  $L$  is a reference length and is defined as  $L = \mu \nu / \gamma \sigma_0 T_0$  and  $\delta^*$  can be introduced as  $\delta^* = (\gamma^* c_0 + \gamma^{**} N_0) / \gamma T_0$ .

physical quantities of practical significance in this work are the local Nusselt number  $Nu$ , and local Sherwood number  $Sh$ , and the local density number of the motile micro-organisms  $N_n$  which are expressed as

$$Nu = -\frac{x \left( \frac{\partial T}{\partial y} \right)_{y=0}}{(T_w - T_\infty)}, \quad Sh = -\frac{x \left( \frac{\partial C}{\partial y} \right)_{y=0}}{(C_w - C_\infty)}, \quad N_n = -\frac{x \left( \frac{\partial N}{\partial y} \right)_{y=0}}{(N_w - N_\infty)},$$

$$C_f = x R_e^{-2} f''(0), \quad Nu = -\frac{x}{L} \theta'(0), \quad Sh = -\frac{x}{L} \phi'(0), \quad N_n = -\frac{x}{L} x'(0)$$

Where  $U_w$  is the uniform stretching velocity of the surface.

## Results and Discussion

In this segment, governing coupled ODEs [7-10] with suitable boundary conditions [11] have been solved numerically by utilizing the Runge-Kutta-Fehlberg method with the help of computational software MATHEMATICA.

Table 1 recognizes the fluctuations of the amended Nusselt number,  $-\theta''(0)$ , the amended Sherwood number,  $-\phi''(0)$  and the amended density number of the motile Microbe  $-\chi''(0)$ . The effects of the magnetic parameter  $\xi$  on the fluid velocity boundary layers have tendency to produce a Lorentz force which acts in the direction opposite so that of the flow, causing a flow retardation effect. This causes the fluid velocity to decrease, as seen in Figure 2(a). On the other hand, enhancing the values of  $\xi$  grows the widths of the thermal, the concentration and density of motile microorganism profiles, as seen in Figures 2(b), 2(c) and 2(d). The relationship between a fluid's viscosity and thermal conductivity is known as the Prandtl number, which has no dimensions. As a result, it evaluates the relationship between a fluid's velocity and

its ability to transmit heat. Recognized that a larger Prandtl number fluid has a drop in heat conduction. For larger value of  $Pr$ , the fluid thermal boundary layer will be thinner, as shown in Figure 3(a). As depicted from Figure 3(b), the fluid concentration is

increased by increasing the value of  $Pr$  close to the surface, but the reverse situation is settled apart from the surface. As the fluid thermal boundary layer behaved, the density of motile microorganism profile behaves as shown in Figure 3(c).

**Table 1:** Fluctuations of amended Nusselt/Sherwood numbers.

| $\xi$ | $Pr$ | $Nb$ | $Sc$ | $\lambda$ | $Sc$ | $Le$ | $Pe$ | $\delta$ | $\delta^*$ | $-\theta'(0)$ | $-\phi'(0)$ | $-x'(0)$ |
|-------|------|------|------|-----------|------|------|------|----------|------------|---------------|-------------|----------|
| 0.1   | 0.71 | 0.5  | 0.5  | 0.1       | 0.5  | 3    | 0.5  | 0.5      | 2          | 1.645715      | 0.453097    | 9.746505 |
| 0.5   |      |      |      |           |      |      |      |          |            | 1.589738      | 0.411041    | 9.513847 |
| 0.9   |      |      |      |           |      |      |      |          |            | 1.535616      | 0.372671    | 9.294786 |
| 0.3   | 0.71 | 0.5  | 0.1  | 0.1       | 0.5  | 1    | 0.5  | 0.5      | 2          | 1.61749       | 0.431624    | 5.422875 |
|       | 1    |      |      |           |      |      |      |          |            | 1.932846      | 0.187554    | 5.348174 |
|       | 2    |      |      |           |      |      |      |          |            | 2.628735      | -0.37477    | 5.180795 |
| 0.3   | 0.71 | 0.5  | 0.5  | 0.5       | 0.5  | 0.5  | 0.5  | 0.5      | 2          | 1.499851      | 0.211612    | 3.745131 |
|       |      | 1    |      |           |      |      |      |          |            | 1.356942      | 1.072332    | 3.883243 |
|       |      | 1.5  |      |           |      |      |      |          |            | 1.230011      | 1.25438     | 3.927358 |
| 0.3   | 0.71 | 0.5  | 0    | 0.5       | 0.5  | 0.5  | 0.5  | 0.5      | 2          | 1.582078      | 1.517237    | 3.993665 |
|       |      |      | 0.2  |           |      |      |      |          |            | 1.547979      | 1.100573    | 3.89299  |
|       |      |      | 0.5  |           |      |      |      |          |            | 1.499851      | 0.511612    | 3.745131 |
| 0.3   | 0.71 | 0.5  | 0.5  | -0.5      | 0.5  | 3    | 0.5  | 0.5      | 2          | 1.75603       | 0.330439    | 9.584359 |
|       |      |      |      | 0         |      |      |      |          |            | 1.642617      | 0.413727    | 9.620554 |
|       |      |      |      | 0.5       |      |      |      |          |            | 1.499851      | 0.511612    | 9.664395 |
| 0.3   | 0.71 | 0.5  | 0.5  | 0.5       | 0.5  | 0.5  | 0.5  | 0.5      | 2          | 1.499851      | 0.511612    | 3.745131 |
|       |      |      |      |           | 0.7  |      |      |          |            | 1.465512      | 0.993882    | 3.867973 |
|       |      |      |      |           | 1    |      |      |          |            | 1.429505      | 1.582873    | 4.004337 |
| 0.3   | 0.71 | 0.5  | 0.5  | 0.5       | 0.5  | 0.1  | 0.5  | 0.5      | 2          | 1.499851      | 0.511612    | 1.43486  |
|       |      |      |      |           |      | 0.5  |      |          |            | 1.499851      | 0.511612    | 3.745131 |
|       |      |      |      |           |      | 1    |      |          |            | 1.499851      | 0.511612    | 5.450837 |
| 0.3   | 0.71 | 0.5  | 0.5  | 0.5       | 0.5  | 0.5  | 0.1  | 0.5      | 2          | 1.499851      | 0.511612    | 1.697351 |
|       |      |      |      |           |      |      | 0.5  |          |            | 1.499851      | 0.511612    | 1.935772 |
|       |      |      |      |           |      |      | 1    |          |            | 1.499851      | 0.511612    | 3.745131 |
| 0.3   | 0.71 | 0.5  | 0.5  | 0.5       | 0.5  | 0.5  | 0.5  | 0.1      | 2          | 1.499851      | 0.511612    | 2.609682 |
|       |      |      |      |           |      |      |      | 0.3      |            | 1.499851      | 0.511612    | 3.051287 |
|       |      |      |      |           |      |      |      | 0.5      |            | 1.499851      | 0.511612    | 3.745131 |
| 0.3   | 0.71 | 0.5  | 0.5  | 0.5       | 0.5  | 0.5  | 0.5  | 0.5      | 0.5        | 1.070815      | 0.472774    | 2.957496 |
|       |      |      |      |           |      |      |      |          | 1          | 1.245194      | 0.481294    | 3.261206 |
|       |      |      |      |           |      |      |      |          | 1.5        | 1.383445      | 0.495637    | 3.519114 |

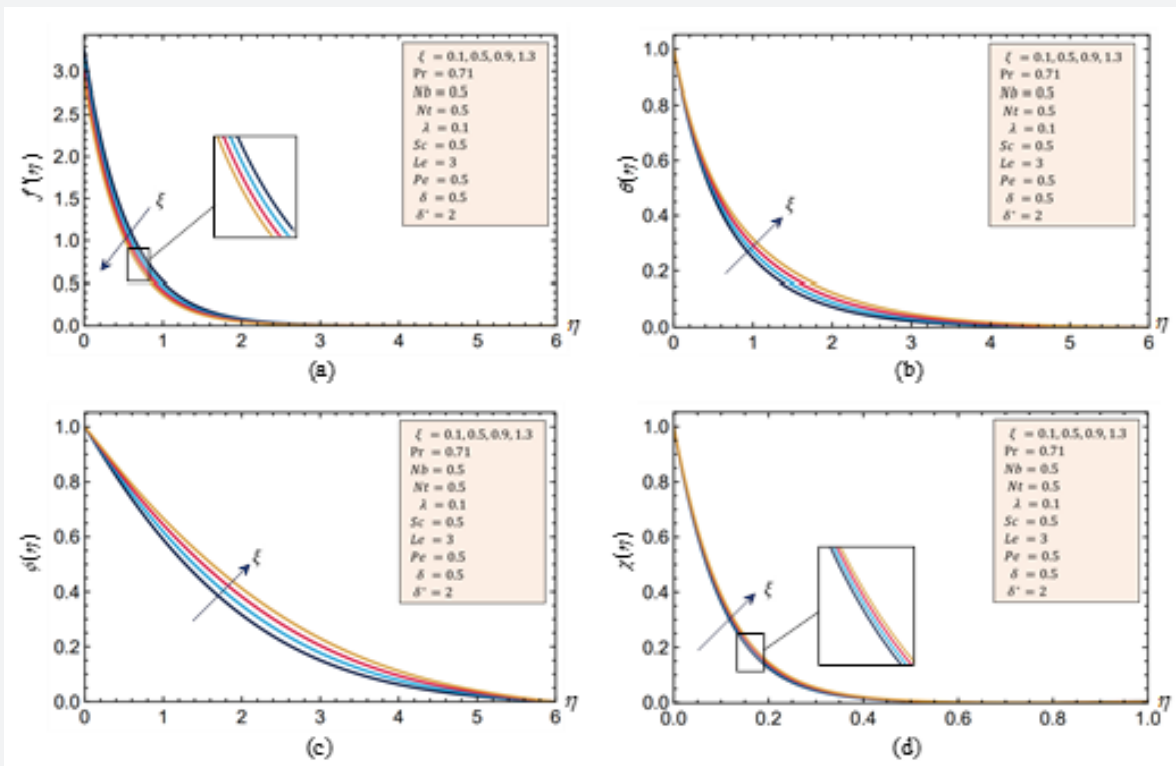


Figure 2: Velocity, temperature, concentration, and motile microorganism density charts influenced by  $\xi$ .

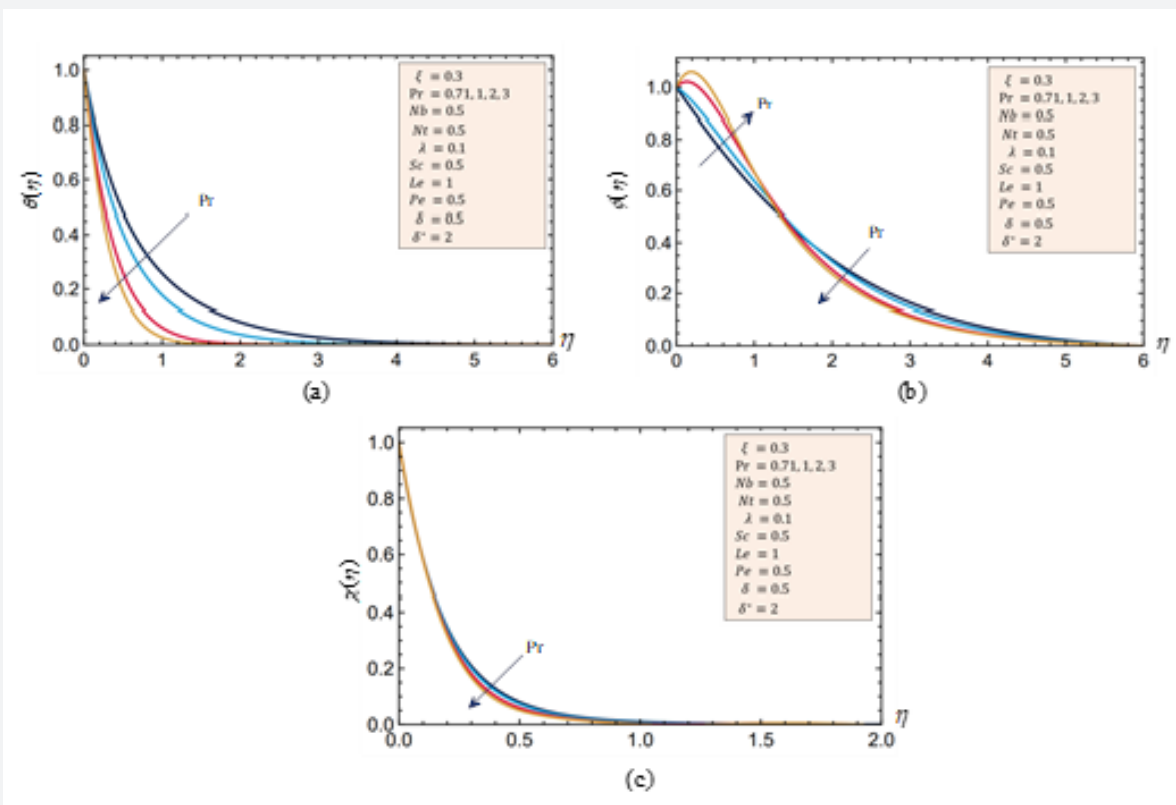


Figure 3: Temperature, concentration, motile microorganism density charts influenced by  $Pr$ .

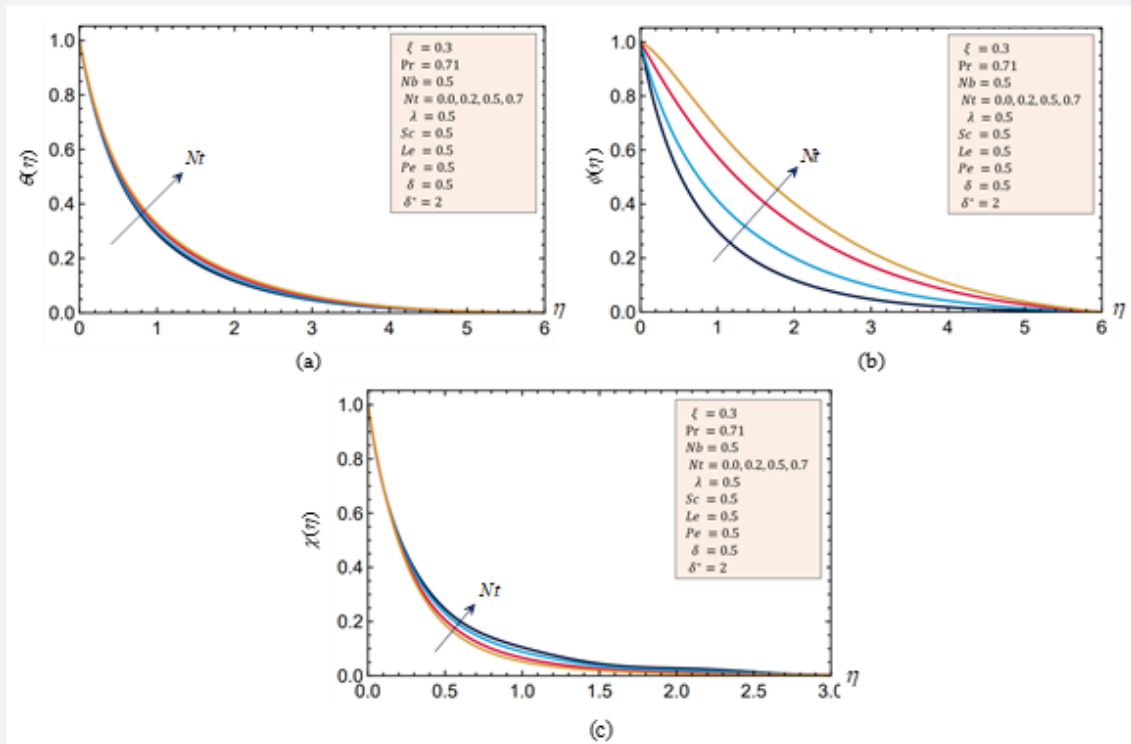


Figure 4: Temperature, concentration, motile microorganism density charts influenced by  $Nr$ .

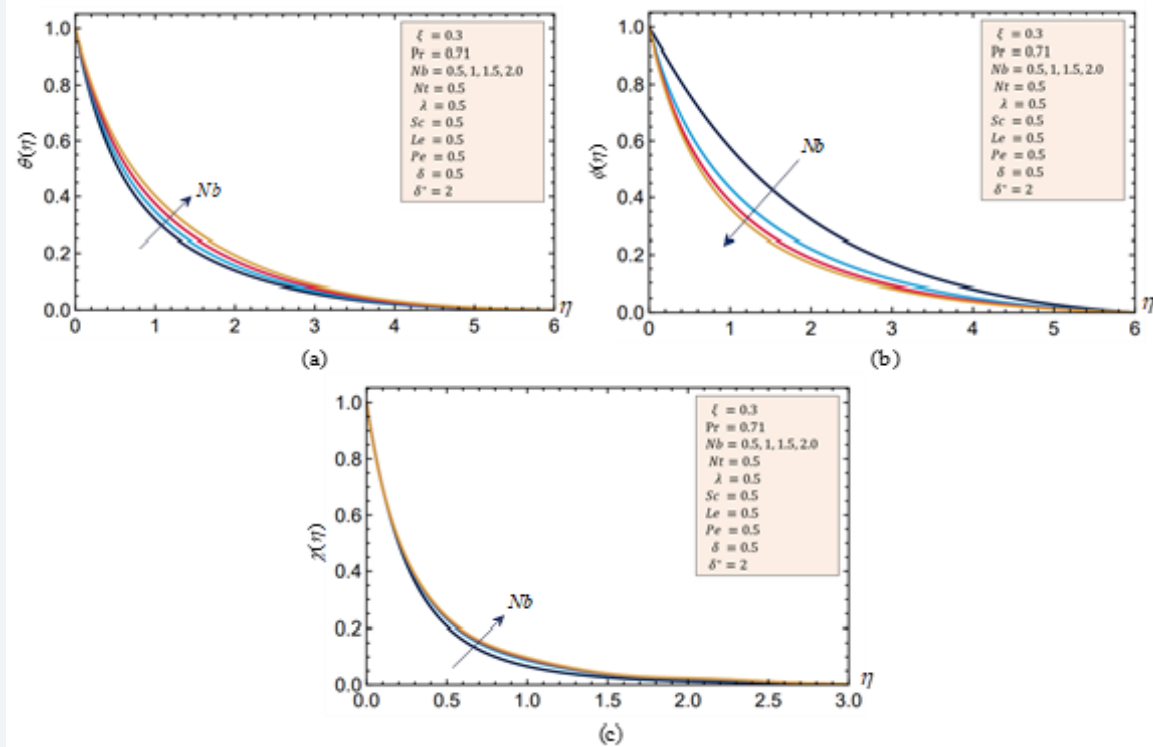


Figure 5: Temperature, concentration, motile microorganism density charts influenced by  $Nb$ .

Thermophoresis is a mechanism where various types of mobile particle mixes respond differently to the force of a temperature gradient. This study uses the thermophoresis parameter  $Nt$  to explain this occurrence. It's important to note that the thermophoresis parameter can have a positive or negative value, with the negative value of  $Nt$  denoting a hot surface and the positive value denoting a cool surface. A hot surface repels particles and motile microbes, as seen in Figures 4(b) and Figures 4(c), therefore thermophoresis tends to blast them away from the surface for hot surfaces. As a result, layer that is largely devoid of nanoparticles and mobile microbes develops just below the surface. As can be seen from Figure 4(a), raising  $Nt$  raises the temperature of the nanofluid. These findings can be explained by the fact that the nanofluid with the lower thermophoresis diffusion parameter has a greater volumetric heat capacity. As a result, the surface cools down and more heat is transmitted from the surface to the nanofluid. The random movement of nanoparticles floating in a fluid

known as Brownian diffusion is the result of their bombardment by the fluid's quickly moving atoms or molecules. The temperature and particle concentration in the boundary layer above the surface are controlled by this motion. The key to understanding Brownian motion is the Brownian motion parameter  $Nb$ . As can be seen from Figure 5(a), increasing  $Nb$  raises the temperature of the nanofluid. The nanofluid with the lowest Brownian diffusion parameter and the maximum volumetric heat capacity is one way to interpret these data. As a result, the surface cools down and more heat is transmitted from the surface to the nanofluid. On the other hand, as can be observed from Figure 5(b), increasing  $Nb$  causes the nanofluid concentration to drop. This effect is explained by the fact that an increase in Brownian diffusion would cause the nanoparticles to scatter more randomly away from the nanofluid's (higher concentration zone) surface (lower concentration zone). The density of the motile Microbe is slightly increased by increasing  $Nb$  as seen from Figure 5(c), (Table 1).

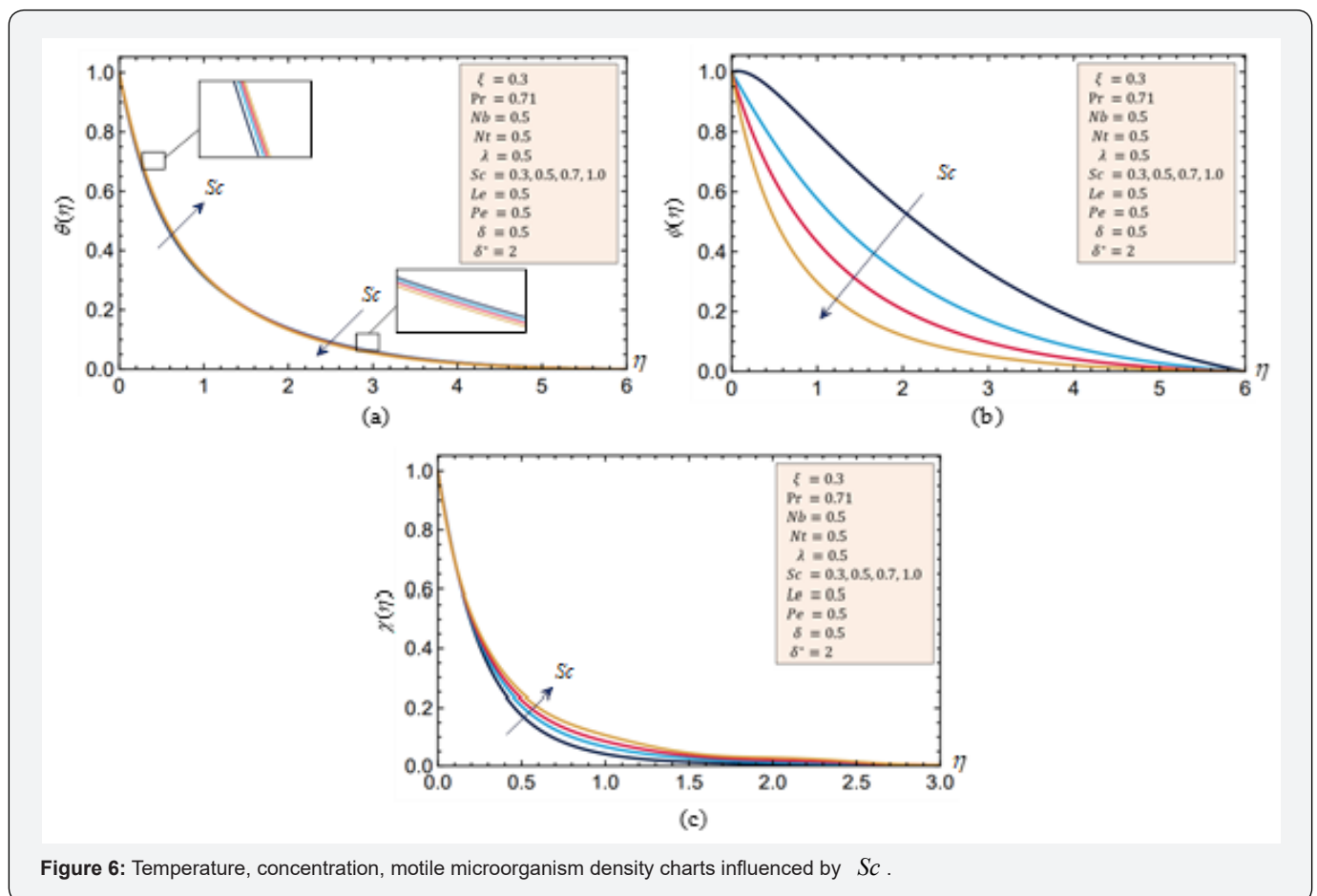


Figure 6: Temperature, concentration, motile microorganism density charts influenced by  $Sc$ .

Figure 6(c) depicts the effect of Schmidt number  $Sc$  on the density of motile Microbe. Close to the surface, the nanofluid temperature is increased slightly by increasing  $Sc$  while the reverse happens apart from the surface as exhibited in Figure 6(a). For bigger values of  $Sc$ , there is a downward tendency of concentration as seen from Figure 6(b). Physically, the concentration and

accompanying layers are reduced due to a decrease in mass diffusivity for bigger  $Sc$ . Heat generation/absorption,  $\lambda$ , is a non-dimensional metric that is dependent on the quantity of heat created or absorbed in the fluid. The fluid temperature is increased as figured out from Figure 7(a) which is anticipated whereas the heat generating engine enhances heat transferred from surface to

fluid and reverse takes place in case of heat absorption. Close to the surface, the nanofluid concentration is decreased slightly by increasing  $\lambda$  while the reverse happens apart from the surface as exhibited in Figure 7(b). The ratio of heat diffusivity to microorganism diffusivity is known as the bio-convection Lewis number. It is used to depict fluid flows where simultaneous microorganism and heat transfer is occurring. Figure 8(a) shows how the bio-convection Lewis number affects the motile microorganism profile.

In the boundary layer regime, bio-convection Lewis number is a measure of how much thermal diffusion rate contributes to microorganism diffusion rate. As a result, when  $Le=1$ , both heat and microorganism diffuse at the same rate, and when  $Le>1$ , heat diffuses more quickly than microorganism. As the bio-convection Lewis number rises, the motile microorganism boundary layer becomes thinner (Figure 2).

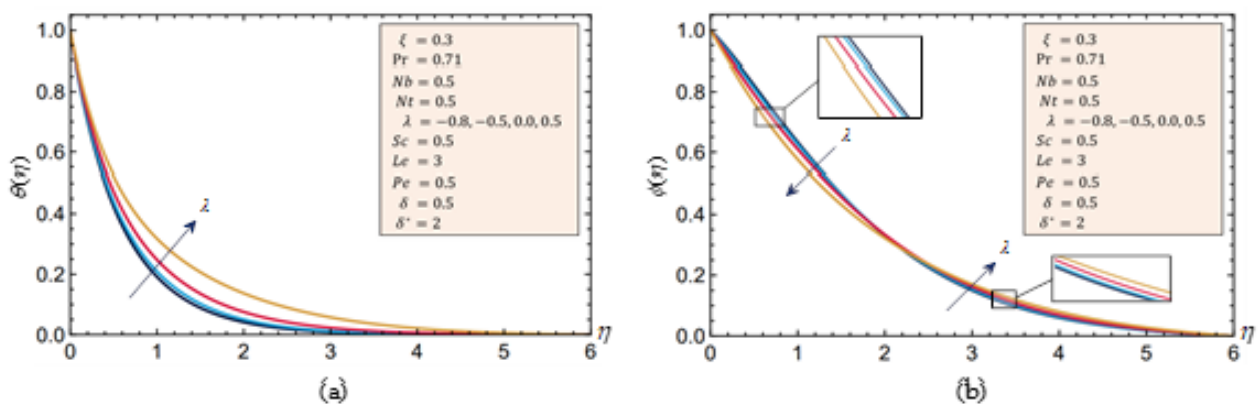


Figure 7: Temperature/concentration charts influenced by  $\lambda$ .

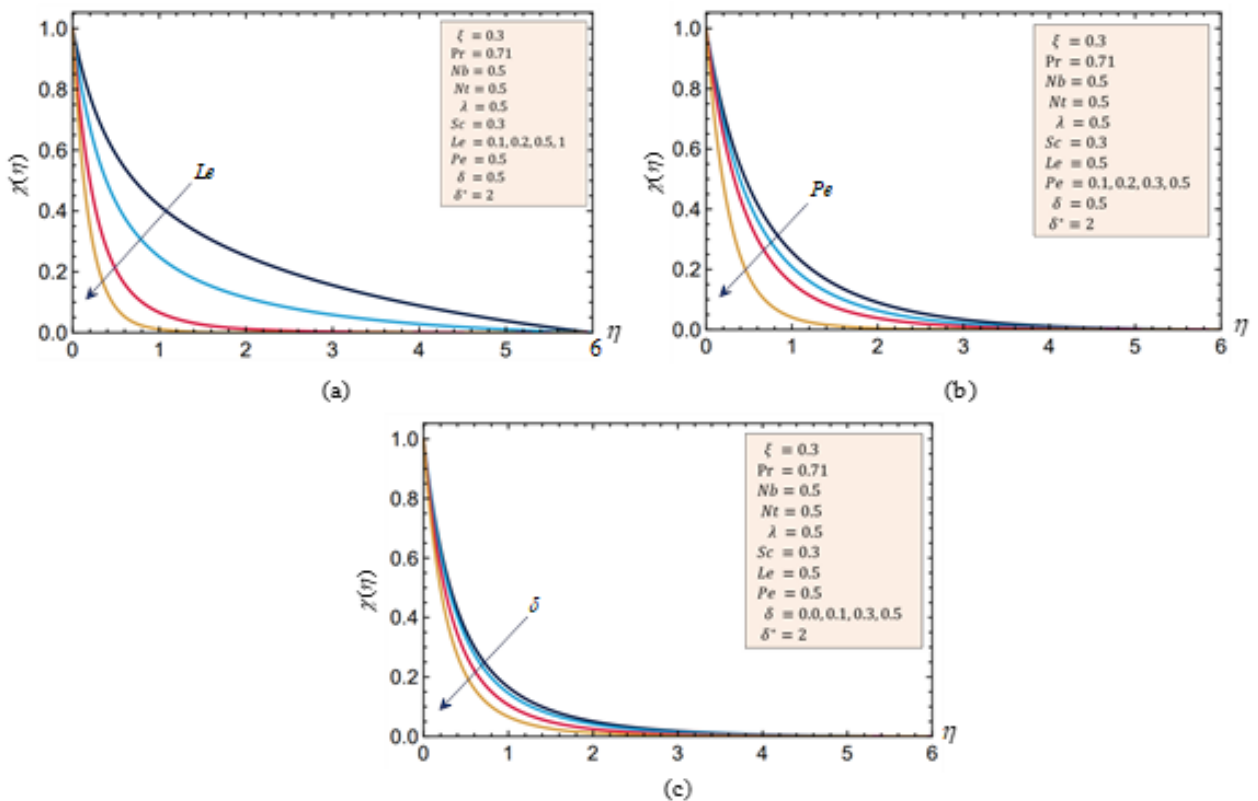
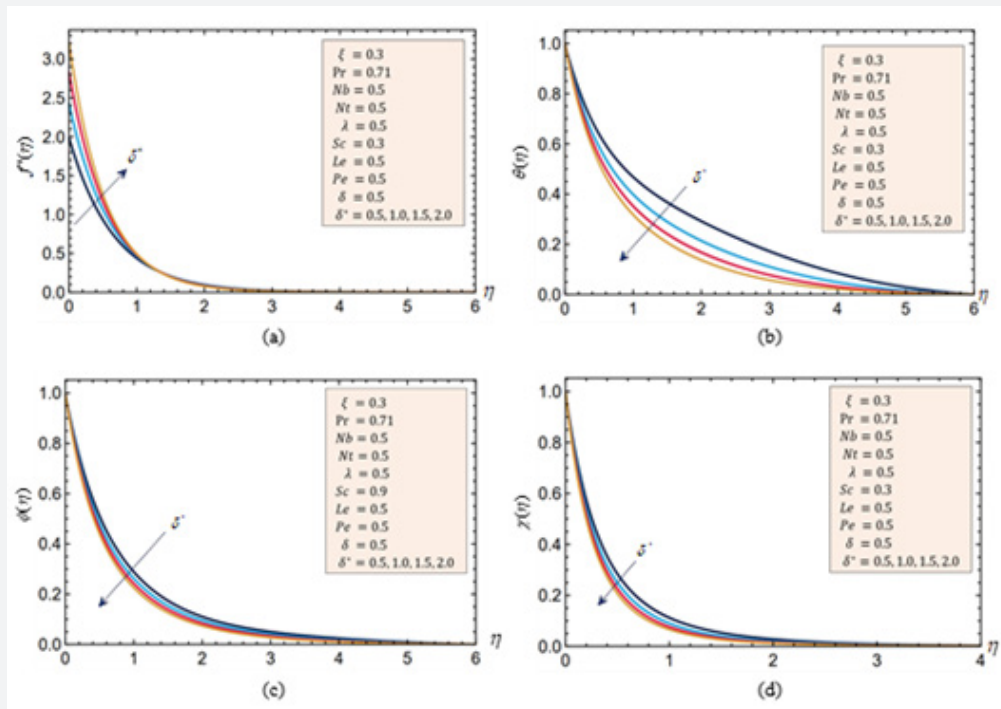


Figure 8: Motile microorganism density charts influenced by  $Le$ ,  $Pe$  and  $\delta$ , respectively.





**Figure 9:** Velocity, temperature, concentration, and motile microorganism density charts influenced by  $\delta^*$ .

A big bio-convection Peclet number suggests an advectively dominated distribution while a small number shows a diffuse flow. The bio-convection Peclet number measures the relative significance of advection vs diffusion. As a result, if the bio-convection Peclet number rises, the motile microorganism boundary layer becomes thinner as shown in Figure 8(b). An increase of motile Microbe parameter  $\delta$  reduces the motile microorganism boundary layer as shown in Figure 8(c) (Figure 3-6). Examining the impact of surface tension on motion in light of the Marangoni convection effect is one of the main objectives of this study. With both heat and mass transmission, the physical values of  $\gamma$  and  $\gamma^*$  represent the gradient of the surface tension. When a result, as the surface tension parameter,  $\delta^*$ , increases, the surface tension decreases (according to the negative values of the coefficients  $\gamma$  and  $\gamma^*$ ). As a result, lowering the surface tension at the free surface lowers the free surface stiffness and enhances the free surface's capacity to move. As a result, as illustrated in Figure 9(a), the velocity increases (Figure 7,8). In addition, Figures 9(b), 9(c) and 9(d) show the impact of surface tension coefficients on heat, concentration, and motile Microbe, respectively. The conditions  $T < T_\infty$ ,  $C < C_\infty$  and  $N < N_\infty$  are considered to be present throughout the boundary layer. In terms of physics, raising both the gradients of temperature, concentration and motile-microorganism density causes more instability at the contact by lowering the surface tension. As illustrated, raising the surface tension parameter lowers the temperature's, the concentrations, and the density of motile-microorganism's profiles.

## Conclusion

The purpose of current investigation is to explore the impacts of heat production on Marangoni flow of a nanofluid containing gyrotactic Microbe across a stretched surface in the presence of magnetic field. Using an appropriate similarity transformation, the controlling boundary layer equations are written down and converted into ordinary differential equations. The numerically solve the resultant ordinary differential equations. Discussion and analysis are conducted on the impacts of the various factors on the velocity, temperature, concentration, and motile microbe density. Additionally, computed and explored for various embedded factors in the issue statements are the density of the motile microbe, the Nusselt number, Sherwood number, and the skin friction coefficient. The findings are visually displayed. It has been discovered that: The fluid velocity is reduced by enhancing magnetic parameter characteristic while it is increased by improving the surface tension. The fluid temperature is enhanced by enhancing permeability, thermophoresis, Brownian motion, or heat generation/absorption characteristic while it is decreased by improving the Prandtl number or the surface tension. Concentration levels are improved with higher permeability parameter or thermophoresis while they are decreased by higher Brownian motion, Schmidt number or surface tension (Figure 9). Density of motile microorganism is improved with higher magnetic parameter, thermophoresis, Brownian motion, or Schmidt number while it is decreased by higher Prandtl number, Lewis number, Peclet number, or surface tension. Close to the surface, the concentration levels are in-

creased by increasing the Prandtl number while the reverse happens apart from the surface. Close to the surface, the temperature profile is slightly increased by increasing the Schmidt number while the reverse takes place apart from the surface. Close to the surface, the concentration levels are slightly decreased by increasing the heat generation/absorption parameter while the reverse happens apart from the surface.

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