



Research Article
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Turbulence and Mixing in Energetic Environment of the Coastal Waters of Sri Lanka



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Abstract

In order to obtain first insights on the ocean turbulence and mixing in the waters around Sri Lanka, the microstructure measurements in the South Monsoon Current to the south of Weligama (along 80.5° E and near the shelf break) and in the East Indian Coastal Current in the Bay of Bengal to the east of Trincomalee (along 8° N) was conducted by National Aquatic Resources Research and Development Agency (NARA) of Sri Lanka aboard the research vessel Samuddrika. It is shown that in strongly startified pycnocline, which is dynamically detached from the low-saline surface mixed layer, the probability distributions of the logarithm of the kinetic energy dissipation rate $log_{in} \varepsilon$ follow generalized extreme value (GEV) model. This could be associated with the occurrence of rare highly-energetic turbulent patches in the pycnocline due to intermittent breaking of internal-wave and sporadic shear-induced instabilities. The mixing Reynolds number (a nomalized eddy diffusivity) was found to be a function of the gradient Richardson number, and followed the parameterization proposed by [1].

Keywords: Bay of Bengal; Indian Ocean; Mixing; Pycnocline; Turbulence

Introduction

The Arabian Sea and Bay of Bengal (BoB) are two major basins in the northern Indian Ocean, which influenced by reversing monsoon wind [2]. Surface salinity of the BoB is greatly affected by large freshwater input due to precipitation (\sim 2 m yr⁻¹) [24] and large river discharges ($1.625 \times 10^{12} \, \mathrm{m^3 yr^{-1}}$) [3]. The Arabian Sea freshwater evaporation is \sim 1 m yr⁻¹. Such imbalance in freshwater fluxes influences thermohaline circulation in the northern Indian Ocean [4]. Large freshwater influx to the northern BoB produces a low-saline shallow mixed layer overlying highly stratified sharp thermohalocline [5]. Understanding small-scale mixing in the pycnocline is crucial for heat, mass, momentum, and biogeochemical fluxes in the Bay [6,7].

Formation of mesoscale eddies and inertial waves [8,9] in the BoB stir water in the interior of the Bay, leading to a cascade of smaller scale local instabilities. Turbulence so generated produces mixing and vertical exchange of physical properties as well as biogeochemical fluxes, influencing the BoB ecosystem and strong near surface stratification in the region restricts turbulence to a thin surface boundary layer [10]. Away from local frontal zones, turbulence in the surface low-salinity layer appears to be decoupled from the underlying pycnocline, where turbulence

occurs only as rare and sporadic events [11,12]. Convergence and divergence of seasonal currents as well as their reversals during monsoon transitions produce remarkable variability of mesoscale and small-scale phenomena near the eastern and southern coastal waters of Sri Lanka [13,14].

Until recently, only a few direct measurements of small-scale turbulence and related mixing in BoB have been conducted [11,15]. As a continuation of this works, in April and September of 2014, Sri Lankan oceanographers launched first ever measurements of microstructure and associated hydrodynamic variables (temperature, salinity, stratification, currents) to the south and to the east from the Sri Lanka coast (Figure 1). The measurements were taken using research vessel R/V Samuddrika of the National Aquatic Resources Research and Development Agency (NARA) under the auspices of the Air-Sea Interactions Regional Initiative (ASIRI) [16].

Oceanographic Measurements Around Sri Lanka

A 60 miles long longitudinal transect (11 stations with Sea-Bird conductivity temperature depth (CTD) and VMP-500 microstructure profiles) were taken by R/V Samuddrika on April

23-25, 2014 along 80.5°E to the south from Weligama (WS). In addition, 14 VMP profiles were obtained during a 3 hrs drift (WD) near the Weligama shelf break. Another 60 miles latitudinal section was completed on September 9 – 12, 2014 along \sim 8°N to

the east of Trincomalee (TS), covering 9 stations with Sea Bird CTD and VMP casts. In addition, ocean currents were measured using a 150 kHz acoustic Doppler current profiler (ADCP) mounted on ship hull.

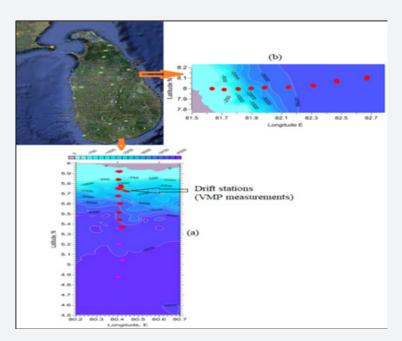


Figure 1: a) R/V Samuddrika measurements taken in April 2014 to the south from the Weligama coast (WS). The transect stations are shown by circles; 14 VMP casts at a drift station (WD) near the shelf break; b) September 2014 section (TS) taken in the BoB to the east from Trincomalee.

The tethered VMP-500 profiler of Rockland Scientific was deployed through the rear deck A-frame down to 150 m depth with almost constant sinking velocity of 0.7 ms $^{-1}$. The VMP carried two airfoil probes to estimate the turbulent kinetic energy [TKE] dissipation rate (\mathcal{E}), an accelerometer, a pressure sensor, and a temperature-conductivity Sea-Bird unit to obtain precise estimates of temperature, salinity, and potential density. Data above first 5 m of the water column were discarded due to potential ship-wake contamination. Data processing followed the methodology of [17].

The TKE dissipation rate ε was calculated by fitting empirical small-scale shear spectra to a Nasmyth benchmark spectrum at consecutive two-second segments (1024 points). As a result, vertical profiles of $\varepsilon(z)$ were obtained with a vertical resolution of ~ 1.4 m. The same spacing was adopted for temperature T(z), salinity S(z), density P(z), and buoyancy frequency P(z) profiles. Here $P(z) = (g/\rho_0) \times (d\rho_0/dz)$ is the squared buoyancy frequency and the corresponding variables are vertical coordinate $Z(z) = (g/\rho_0) \times (d\rho_0/dz)$, potential $P(z) = (g/\rho_0) \times (d\rho_0/dz)^2 + (g/\rho_0/\Delta z)^2 + (g/\rho_0/\Delta z)^2$

Results and Discussion

Stratification to the south and to the east from Sri Lanka During the inter-monsoon season (April 2014), the mixed layer depth (MLD) along the Weligama section did not change much, varying between approximately 20 and 30 m (Figure 2). Surface mixed layer was underlined by strong narrow density jump and then density gradually increases with depth in the main pycnocline below $z\sim\!150\,$ m. Interestingly to noticed that the turbulence measurements carried out in early February of 2014 at WS, only one station of near the shelf break, showed that the MLD was $\sim\!60\,$ m, indicating the possibility of substantial convective cooling and/or strong wind mixing in the upper layer south of Sri Lanka during the winter monsoon.

In September 2014, at the end of the southern monsoon season, stratification in the upper 150 m along the Trincomalee section (TS) was much more complex compared to that observed along Weligama (WS) in April. The density structure in Figure 2b exhibits a sharp pycnocline in the depth range z=30-40 m not far from the shelf break. About 50 miles offshore, however, the pycnocline upwelled towards the sea surface, forming a striking baroclinic (density) front that separates diluted BoB surface water,

which is moving southward along the Sri Lanka coast, from warm, saltier water at the eastern end of the transect [16]. The latter represents the Arabian Sea water moving north/northeastward

as the southwest monsoon current [18]. The shipboard ADCP measurements support this concept [8].

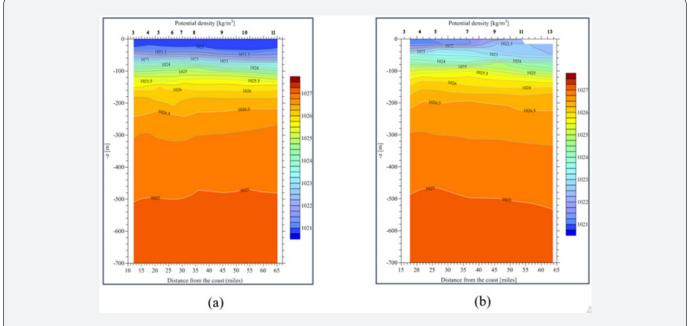


Figure 2: Potential density contour plots (Seabird CTD data): a) the southern section (WS, April 2014) and b) the eastern section (TS, September, 2014).

The Dissipation Rate Statistics

It has been suggested [19] that the probability distribution of the logarithm of the TKE dissipation rate $\log_{10} \varepsilon$ in stably stratified pycnocline may follow generalized extreme value distribution [20] due to rare appearance of energetic turbulence events. Accordingly, a small number of patches with high level of TKE dissipation rate exist, while most of the background turbulence is weakly dissipative or consists of smaller patches in their final stages of life. The cumulative distribution functions (*CDF*) of the generalized extreme value (GEVD) model are expressed as Eq. 1a.

$$CDF\left(\tilde{x}\right) = \exp\left\{-\left[1 + \xi \tilde{x}\right]^{-1/\xi}\right\} \tag{1a}$$

$$1 + \xi \tilde{x} > 0 \qquad \text{and} \qquad \xi \neq 0 \tag{1b}$$

where $\tilde{x} = \left(x - \mu_{gev}\right) / \sigma_{gev}$ (in our case $x \equiv \log_{10} \varepsilon$) and ξ , σ_{gev} , and μ_{gev} are the shape, scale, and location parameters, respectively [19].

We calculated ${\it CDF}(\varepsilon)$ for the dissipation rate data along WS and TS shown in Figs. 4 a, b of [11] along a mini transect WD of 14 drift stations. The dissipation rate at the WD is shown in Figure 3. The empirical cumulative distribution functions ${\it CDF}(\varepsilon)$ are shown in Figure 4 separately for the southern (WS) and eastern (TS) transects and for drift measurement (WD). The GEVD approximates all three empirical distributions well. The scale

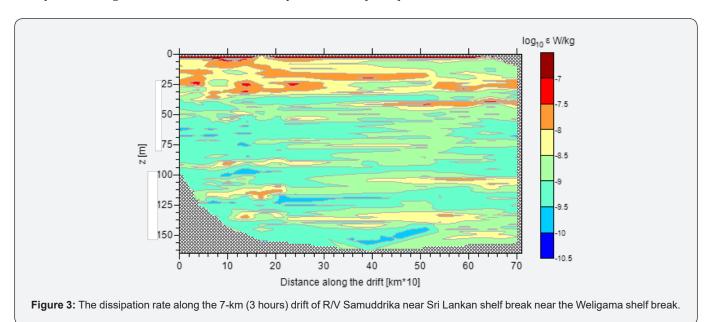
 $\sigma \sim (0.28-0.43)$ and location $\mu \sim -(9.1-9.6)$ parameters are quite similar for all regions, while the shape parameter ξ can be positive or negative $\sim \pm 0.07$. For ξ close to zero, the probability distribution function is almost symmetric, and it is highly skewed towards small values of \tilde{x} , if ξ is negative (e.g., see WS).

The empirical distributions of $\log_{10} \varepsilon$ shown in Figure 4 were also compared with a more traditional normal model using the Akaike Information Criterion AIC [21]. It appears that for all analyzed distributions, the AIC is higher for the GEVD model than for the normal model. The difference between the AIC values is statically significant. On the contrary to the observations of pycnocline turbulence, the probability distribution functions $\log_{10} \varepsilon$ for the very weakly stratified, almost homogeneous, upper layer (the depth range between MLD and z=10 m) generally follow the normal distribution (not shown here for brevity). Log-normal distribution for the TKE dissipation rate has been formulated by [22] for homogeneous locally isotropic turbulence and has been well supported by numerous experimental data [19,23,24] whereas $\log_{10} \varepsilon$ in marine and lacustrine quasi-homogeneous boundary layers often follow the normal distribution.

From the probabilistic point of view, the generation/dissipation of intermittent intense energetic turbulence in the strongly startified pycnocline can be considered as a random sequence of rare events. The sources of such turbulence are most probably associated with non-stationary, intermittent internal-

wave breaking and sporadic shear-induced instabilities [6]. In weakly stratified regions, sub-mesoscale intermittency of the TKE

dissipation rate follows more traditional log-normal distribution [23.25].



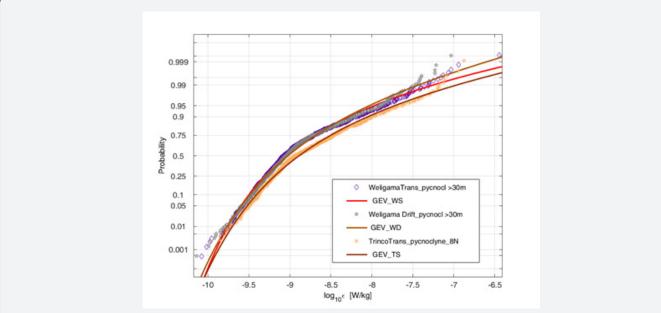


Figure 4: The cumulative probability distributions of the logarithm of the TKE dissipation rate in the pycnocline for WS, WD and TS, approximated by generalized extreme value (GEV) distributions. The inset indicates the depth below the upper boundary of the pycnocline for each dataset. The pycnocline is of variable depths along TS.

Eddy Diffusivity vs The Richardson Number

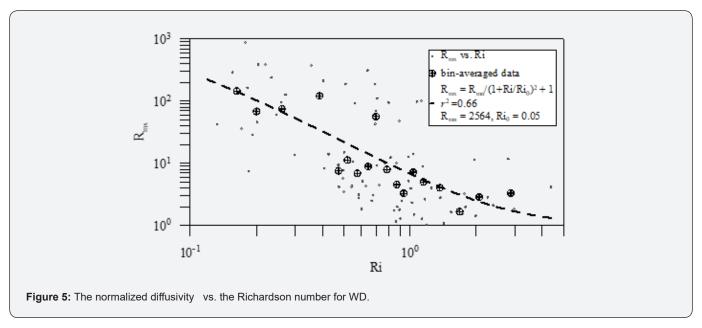
Jinadasa et al. [11] suggested that shear-induced turbulence is a major possible source of diapycnal mixing in monsoon driven currents around Sri Lanka, and if so, the gradient Richardson number is a salient parameter that determines the eddy diffusivity. To verify this hypothesis, we analyzed the dependence of turbulent diffusivity, $K_N = 0.2 \, \varepsilon/N^2$ on the gradient Richardson

number as in many previous publications [26-30]. We used 14 VMP profiles collected during the 3 hr (7 km) drift of R/V Samuddrika along the Weligama shelf break. The dissipation field along this drift is shown in Figure 3 and appears to be typical of stratification in waters all around Sri Lanka away from the frontal zones as reported by [11].

The estimated normalized turbulent diffusivity or the mixing Reynolds number $R_{mx} = K_N/\nu_T$, where $K_N = 0.2\,\varepsilon/N^2$ and ν_T are the eddy and turbulent diffusivities, respectively, is shown in Figure 5 as a function of the gradient Richardson number (Ri). The binaveraged (5 points per bin) samples of both variables were found to minimize scatter. Although the scatter of bin-averaged samples

in Figure 5 is still relatively high (the coefficient of determination of the approximation (1) $r^2 = 0.66$ is less than 0.8), the general decreasing trend of R_{mx} with growing Ri is evident. The dashed line in Figure 5 approximates $K_N(Ri)$ by formula used by [23].

$$R_{mx} = \frac{R_{mn}}{(1 + Ri/Ri_0)^s} + R_{mb}$$
 (2)



where turbulence in neutrally stratified flow is specified by the constant $R_{mn}=2564$ (the corresponding diffusivity $K_N\sim 5.1\times 10^{-4}~{\rm m}^2/s$), $Ri_0=0.05$ and $R_{mb}=1$ corresponds to background molecular mixing. The values of Ri_0 and s=2 are used for fitting based on [23], although the latter authors analyzed atmospheric stably stratified turbulence. The agreement indicates the robustness of (2) for mixing in strongly stratified natural shear flows, and that thermocline mixing around Sri Lanka indeed can be caused by shear instability [31-33].

Conclusion

It is found that statistics of the kinetic energy dissipation rate in the oceanic *pycnocline* around Sri Lanka (TS and WS transects) can be well defined by generalized extreme value distribution (GEVD), which approximates the empirical cumulative distribution functions of the logarithm of the TKE dissipation rate $log_{10} \, \varepsilon$, supporting findings of [19] in the east China Sea In weakly stratified almost homogeneous surface layer (in our case, the depth range between the MLD and z=10 m), the probability function of $log_{10} \, \varepsilon$ generally follows traditional [22] normal distribution. From the probabilistic point of view, the generation/dissipation of intense energetic turbulence in strongly startified pycnocline, which is largely detached from the surface layer of low salinity waters all the way around Sri Lanka, can be considered as randomly occuring rare events.

The sources of such turbulence most probably is associated with non-stationary, shear-induced instabilities induced by vertical shear (e.g., internal wave shear). Based on the measurements taken near the Weligama shelf break, we found that turbulent diffusivity (or mixing Reynolds numebr Re_{mx}) in the pycnocline is a function of the gradient Richardson number Ri (Eq. 2) with fitting parameters consistent with those proposed by [1]. This suggests the robustness of the parameterization as well as the dominance of shear-induced mixing in the thermocline along the trasects analyzed in this study.

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